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## H<sub>2</sub> Dissociation on Noble Metal Single Atom Catalysts Adsorbed on and Doped into CeO<sub>2</sub> (111)

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# H<sub>2</sub> Dissociation on Noble Metal Single Atom Catalysts Adsorbed on and Doped into CeO<sub>2</sub> (111)

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## ABSTRACT:

We used density functional theory (DFT) calculations to investigate the dissociation of H<sub>2</sub> on an Ag single atom catalyst adsorbed on the pristine CeO<sub>2</sub> (111) surface (Ag/CeO<sub>2</sub>), or substituting a surface Ce atom on the reduced (Ag:CeO<sub>2-x</sub>) and partially hydrogenated (Ag:H-CeO<sub>2</sub>) surfaces. The initial state of the H<sub>2</sub> dissociation reaction in the different investigated models corresponds to distinct oxidation states, +1, +2, or +3, of the Ag atom, thus allowing us to examine the influence of the charge transfers between the noble metal, the oxide and the hydrogen atoms on the reaction pathway and activation energy. In all investigated models, the computed barrier of H<sub>2</sub> dissociation is lowered by about 0.6 eV in comparison to that on metal-free CeO<sub>2</sub>. On Ag/CeO<sub>2</sub> and Ag:CeO<sub>2-x</sub>, also the energy of H<sub>2</sub> dissociative adsorption is smaller than on metal-free ceria. These results suggest that CeO<sub>2</sub> modified with dispersed Ag atoms is a promising anode material for proton exchange membrane fuel cells. Further comparison of our results for Ag to analogous calculations for Cu and Au single atom catalysts reveals trends in the computed barriers that can be related to the change of the metal oxidation state in the reaction.

## I. Introduction

Catalysts are essential in many industrial and energy conversion processes. Noble metals have shown excellent catalytic properties for energy production and conversion<sup>1</sup>, but their use on the large scale is limited by the high cost. An important goal is thus to reduce the amount of noble metal while maintaining or possibly increasing the performance. It was recently proposed that this goal could be achieved by single noble metal atoms on a supporting surface<sup>2</sup>. In this so-called single atom catalysis (SAC), the main catalytic center is an isolated supported metal atom. SAC has received increasing attention over the last years, and new techniques to prepare and characterize single atom catalysts have been developed, including wet chemistry, mass selected soft landing, and atomic layer deposition<sup>1</sup>. While such single atom catalysts are expected to exhibit novel properties in comparison to those of nanoparticle catalysts, understanding of their activity is still limited<sup>3</sup>.

In this work, we investigate the reactivity of single noble metal atoms focusing on the dissociation and oxidation of H<sub>2</sub> on Ag, Au and Cu atoms supported on CeO<sub>2</sub> (111), the most stable<sup>4 5 6 7</sup> and frequently exposed surface of CeO<sub>2</sub> nanocrystals. As the reaction occurring at the anodes of proton exchange membrane fuel cells (PEMFCs), the oxidation of H<sub>2</sub> is of both scientific and technological interest. Currently, platinum is the metal of choice for the anode catalyst in PEMFCs but its high cost and scarcity limits its use in industrial applications<sup>8</sup>. Replacing platinum with a cheaper material is crucial for increasing the competitive advantage of fuel cells. A good candidate is cerium oxide, a strongly reducible oxide catalyst<sup>9 11 12</sup> that is also known to be a good electrode material for electrochemical devices<sup>10</sup> and to facilitate the oxidation of H<sub>2</sub>, especially in the presence of noble or transition metals<sup>13</sup>.

As shown by recent experiments, metal atoms on CeO<sub>2</sub> often occupy Ce substitutional sites at low levels of loading<sup>1</sup>, and surface cation (Ce) vacancies have been reported to be essential for stabilizing single atom catalysts<sup>1</sup>. Qiao et al.<sup>14</sup> synthesized single Au atom catalysts supported on CeO<sub>2</sub>, and observed that the Au atoms were indeed accommodated in the Ce vacancies. Based on the

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3 classification of McFarland et al.<sup>15</sup>, Ag, Au, and Cu atoms belong to the class of low valence dopants  
4 (LVDs) for CeO<sub>2</sub>. The main effect of LVDs is to determine an increase of the oxidative power of the  
5 oxygen atoms near the dopant<sup>15</sup>. Consequently, these dopant metal atoms should be good candidates  
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7 for decreasing the activation energy for H<sub>2</sub> dissociation.  
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12 Motivated by the above findings and predictions<sup>14, 15</sup>, in this work we use first principles  
13 calculations to investigate the mechanism of H<sub>2</sub> oxidation on single noble metal atoms either adsorbed  
14 on or doped into CeO<sub>2</sub> (111). We mainly focus on single Ag atom catalysts, and compare the  
15 adsorption and reactivity of Ag to those of Cu and Au, for which several SAC studies are already  
16 available<sup>14, 16</sup>. Three different configurations of the single metal atom (M) catalyst are examined: a)  
17 M adsorbed on the pristine CeO<sub>2</sub> (111) surface; b) M substituting a Ce atom in the presence of an  
18 excess electron supplied by an adsorbed H atom; c) M substituting a Ce atom in the presence of a  
19 nearby oxygen vacancy. These configurations correspond to three different initial oxidation states of  
20 the noble metal atom, +1 for case a), +3 for case b), and +2 for case c), which allows us to obtain  
21 insight into the influence of the metal oxidation state on the activation of H<sub>2</sub>. Oxidation states are  
22 here evaluated through established procedures based on analysis of geometries, atomic charges and  
23 partial densities of states (see Sec. II). Although approximate, this approach is widely used and found  
24 to provide meaningful trends for a wide range of systems. While more accurate methods for  
25 determining oxidation states are now available<sup>17, 18</sup>, they are typically quite demanding and outside  
26 the scope of the present investigation.  
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## 48 **II. Computational Methods**

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50 Calculations were performed using spin-polarized density functional theory (DFT) with the  
51 gradient-corrected Perdew- Burke-Ernzerhof (PBE)<sup>19</sup> exchange and correlation functional as  
52 implemented in the Quantum Espresso Package<sup>20, 21</sup>. The interaction between electrons and ions was  
53 described by ultra-soft pseudopotentials. The wave function and charge density energy cutoffs were  
54 set to 30 Ry and 320 Ry, respectively. To account for electron correlations in CeO<sub>2</sub>, we applied the  
55 Hubbard correction (DFT+U) on the f orbitals of the cerium atoms using the implementation of  
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Cococcioni et al.<sup>22</sup>. The U parameter was set to 4 eV, a value that reproduces experimental and theoretical results of previous studies.

CeO<sub>2</sub> crystallizes in the fluorite structure. The bulk primitive cell was optimized using a 7 × 7 × 7 Monkhorst-Pack (MP) mesh grid<sup>23</sup>; the resulting lattice parameter is a=5.52 Å, in good agreement with the experimental value a<sub>exp</sub>= 5.41 Å<sup>24</sup>. The CeO<sub>2</sub> (111) surface was modeled using slabs with three O-Ce-O trilayers and 15 Å of vacuum to avoid interactions between the repeated images. A (3×3) surface supercell was used, and the irreducible part of the surface Brillouin zone was sampled using a (2×2×1) MP mesh grid<sup>23</sup>. A (4×4) surface supercell was used to study H<sub>2</sub> dissociation on an Ag<sub>2</sub> dimer. Atomic positions were relaxed with convergence thresholds of 10<sup>-5</sup> eV and 1×10<sup>-2</sup> eV/Å for the total energy and ionic forces, respectively. Atoms of the last trilayer were kept fixed in their bulk positions.

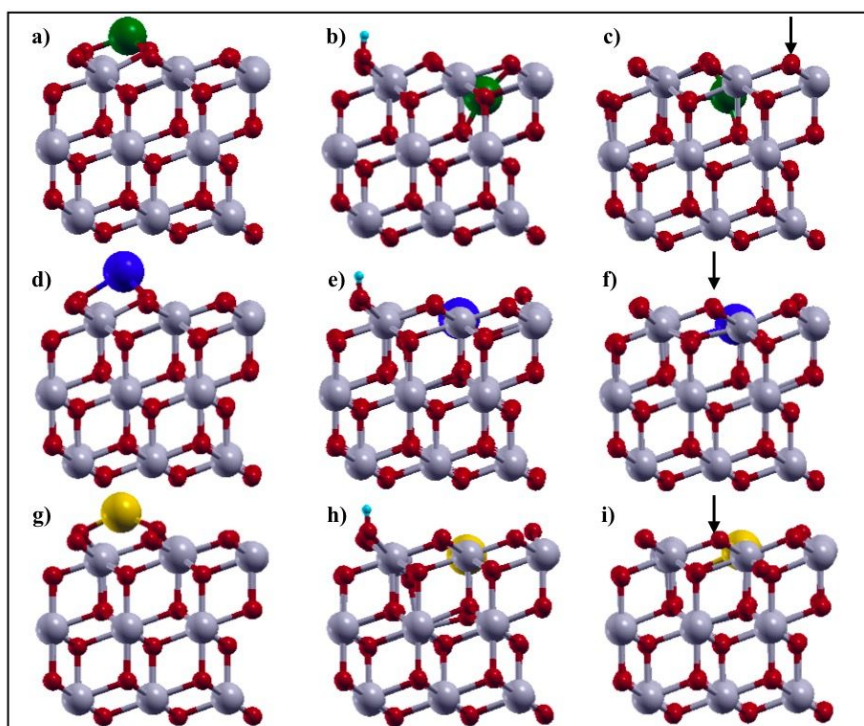
The formation energy of an oxygen vacancy was calculated as:  $E_{\text{form}}=E_{\text{def}}+1/2E_{\text{O}_2}-E_{\text{pristine}}$ , where  $E_{\text{def}}$  is the energy of the surface with the oxygen vacancy,  $E_{\text{pristine}}$  is the energy of the slab without the oxygen vacancy, and  $E_{\text{O}_2}$  is the energy of the O<sub>2</sub> molecule. The adsorption energy of a noble metal atom on the CeO<sub>2</sub> surface was calculated as  $E_{\text{ads}}^M = E_{M/\text{CeO}_2} - E_M - E_{\text{CeO}_2}$ , where  $E_{M/\text{CeO}_2}$  is the energy of the surface with the metal absorbed,  $E_M$  is the energy of the metal in the gas phase, and  $E_{\text{CeO}_2}$  is the energy of the pristine surface. The H<sub>2</sub> adsorption energy was calculated as  $E_{\text{ads}}(\text{H}_2) = E_{\text{H}_2/M@\text{CeO}_2} - E_{\text{H}_2} - E_{M@\text{CeO}_2}$ , where  $E_{\text{H}_2/M@\text{CeO}_2}$  is the energy of the slab with the supported metal atom (M) and the H<sub>2</sub> molecule,  $E_{M@\text{CeO}_2}$  is the energy of the slab with the supported metal atom but without H<sub>2</sub>, and  $E_{\text{H}_2}$  is the energy of the H<sub>2</sub> molecule.  $E_{\text{H}_2}$  was determined using a cubic box with a 20 Å edge; the computed bond length is 0.75 Å, in good agreement with the experimental value (0.740 Å) and other DFT calculations (0.743 Å<sup>16</sup>).

H<sub>2</sub> dissociation pathways were determined using the climbing image nudged method (CI-NEB)<sup>25</sup> with seven images. Charge transfers along the pathways were estimated by analyzing Bader charges<sup>26</sup> and projected density of states (PDOS).

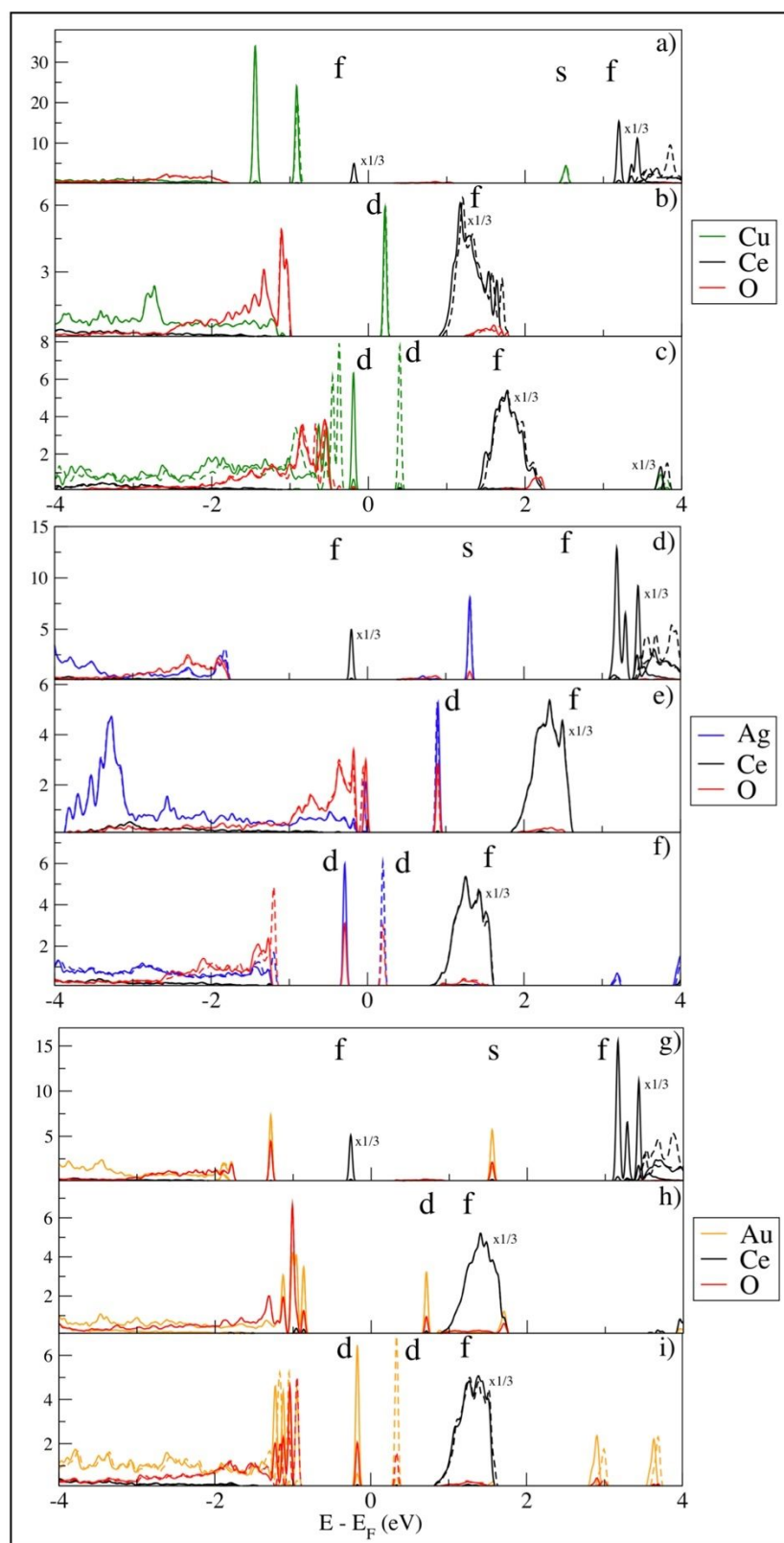
### III. Results and Discussion

#### III.1. Geometries and Electronic Properties of Single Noble Metal Atom Catalysts on CeO<sub>2</sub>(111)

For each single noble metal atom (M), we have studied three different configurations on the CeO<sub>2</sub>(111) surface: an adsorbed M (M/CeO<sub>2</sub>); M substituting a Ce atom in the presence of an excess electron supplied by an adsorbed H atom (M:H-CeO<sub>2</sub>); and M substituting a Ce atom in the presence of a surface oxygen vacancy (M:CeO<sub>2-x</sub>). Fig. 1 and Fig. S1 show side and top views, respectively, of the three configurations while Fig. 2 reports the corresponding orbital projected density of states (PDOS). For each case, we first discuss the results for Ag and next compare the behavior of Ag to the other noble metals (Cu and Au).



**Fig. 1** Models (side views) of a supported noble metal atom M (M = Cu, Ag, Au) on CeO<sub>2</sub>(111) investigated in this work: a), d), g) M adsorbed on CeO<sub>2</sub> (M/CeO<sub>2</sub>); b), e), h) M substituting a Ce atom in the presence of an adsorbed H (M:H-CeO<sub>2</sub>); c), f), i) M substituting a Ce atom in the presence of an Oxygen vacancy (M:CeO<sub>2-x</sub>). Green (first row), blue (2<sup>nd</sup> row), and yellow (third row) spheres indicate the Cu, Ag, and Au atoms, respectively. Gray and red balls represent Cerium and Oxygen atoms. Arrows indicate the positions of the Oxygen atoms that were removed. Top views of these models are shown in Figure S1.



**Fig. 2** Projected Density of States of the model systems shown in Fig. 1: a) Cu/  $\text{CeO}_2$ ; b) Cu:H- $\text{CeO}_2$ ; c) Cu:  $\text{CeO}_{2-x}$ ; d) Ag/  $\text{CeO}_2$ ; e) Ag:H- $\text{CeO}_2$ ; f) Ag:  $\text{CeO}_{2-x}$ ; g) Au/  $\text{CeO}_2$ ; h) Au:H- $\text{CeO}_2$ ; i) Au:  $\text{CeO}_{2-x}$ . Full and dashed lines refer to spin up and spin down contributions, respectively. f states have been multiplied by a factor 1/3.

### III.1.1. *M/CeO<sub>2</sub>*

Fig. 1d shows the structure of an adsorbed Ag atom on CeO<sub>2</sub>(111). Ag binds on top of an oxygen atom of the third layer and is coordinated to three surface oxygen atoms, with adsorption energy of -1.31 eV. The adsorption causes a charge transfer from Ag to the oxide, which results in the formation of a cation Ag<sup>+</sup>, and the reduction of one second neighbor cerium atom from Ce<sup>4+</sup> to Ce<sup>3+</sup> (see Fig. S1), as theoretically found also by Tang et al.<sup>27</sup> Table 1 reports the calculated adsorption energy  $E_{\text{ads}}^{\text{Ag}}$ , the three metal-oxygen bond lengths, the metal – Ce<sup>3+</sup> bond length, the transferred charge, and the magnetization of the reduced Ce<sup>3+</sup> ion, resulting from the partial filling of the *f* orbital. In Fig. 2d we show the projected density of states of Ag/CeO<sub>2</sub>: the appearance of the Ag *s* peak above the Fermi energy ( $E_{\text{F}}$ ) and a new Ce *f* peak below  $E_{\text{F}}$  confirms the charge transfer. The new oxidation state of the Ag adatom is +1, with Ag-O distances slightly larger than the experimental values for Ag<sub>2</sub>O<sup>28</sup>.

The behaviors of Cu and Au are similar to that of Ag. Cu binds on top of a surface oxygen atom of the third layer (Fig. 1d), with a larger adsorption energy ( $E_{\text{ads}}^{\text{Cu}}=-2.62$  eV) in comparison to Ag, and Cu-O distances slightly larger than those observed in Cu<sub>2</sub>O<sup>29</sup>. In contrast to Ag and Cu, Au prefers to adsorb on a bridge site between two surface oxygen atoms, but the adsorption energy difference between the bridge and on top sites is less than 0.10 eV. The Au adsorption energy is  $E_{\text{ads}}^{\text{Au}}=-1.12$  eV. Similar to Ag, Cu and Au are oxidized to Cu<sup>+</sup> and Au<sup>+</sup> upon adsorption on the CeO<sub>2</sub> surface, and at the same time one Ce atom is reduced to Ce<sup>3+</sup>. As shown in Fig. 2a and g, the metal *s* peak is indeed empty, and a new Ce *f* peak appears below  $E_{\text{F}}$ . The reduced Ce atom is a second neighbor of both Cu<sup>27</sup> and Ag, whereas it is a third neighbor in the case of Au (see Table 1), as found in other computational studies<sup>27 30</sup>.

**Table 1** Adsorption energy of the metal adatom, in eV, magnetic moment of  $Ce^{3+}$  ( $\mu_B$ ), M-O distance ( $\text{\AA}$ ) between the metal (M) and neighboring O atoms, M- $Ce^{3+}$  distance ( $\text{\AA}$ ) between the metal and the  $Ce^{3+}$  ion, charge transfer  $\Delta q$  (e) from the metal adatom to the surface.

M	$E_{ads}$	$\mu(Ce^{3+})$	M-O	M- $Ce^{3+}$	$\Delta q$
Cu	-2.62	0.89	2.02, 2.03(2)	4.82	0.55
Ag	-1.31	0.89	2.35, 2.41(2)	5.06	0.38
Au	-1.12	0.90	2.15, 2.18	4.71	0.20

### III.1.2. *M:H-CeO<sub>2</sub>*

The second investigated model is the substitution of a Ce by an Ag atom. Cerium has oxidation state +4 in  $CeO_2$ , whereas the highest oxidation state that Ag can have is +3. To compensate for the missing charge, we have introduced an adsorbed H atom on a distant surface oxygen (see Figs. 1 and S1), which provides an additional electron without substantially affecting the interaction of Ag with the  $H_2$  molecule. We found indeed that, in the absence of such a compensating charge, the  $H_2$  molecule tends to react with an oxygen atom of the  $CeO_2$  surface to form an  $H_2O$  molecule together with a surface oxygen vacancy, suggesting that the surface needs extra electronic charge to be stabilized. The optimized Ag:H- $CeO_2$  structure is shown in Fig. 1e. Ag is approximately at the same height of the Ce surface atoms, and the Ag-O bonds are  $\sim 2.19 \text{\AA}$ , slightly larger than those observed in  $Ag_2O_3$ <sup>31</sup>. The three oxygen atoms bonded to the Ag metal, and the other oxygen atoms below, show a smaller amount of electronic charge ( $\sim -0.2$  e per oxygen atom) than in pure ceria, as found also in a previous computational study for Au<sup>32</sup>. Comparison of the PDOS of Ag:H- $CeO_2$  (Fig. 2e) and Ag/ $CeO_2$  (Fig. 2d) shows a new Ag peak above  $E_F$  for Ag:H- $CeO_2$ , which originates from empty d metal orbitals with some contributions from the *p* states of the less charged oxygen atoms. At the same time, the Bader and Löwdin charges show a depletion of both the *s* and *d* Ag orbitals with respect to the isolated atom. These findings suggest that the new Ag oxidation state is +3, as previously proposed in a theoretical study by Camellone et al.<sup>30</sup> for substitutional Au. This

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3 conclusion is also supported by the Ag - O bond lengths and by the value of the dihedral angle in  
4 Ag:H-CeO<sub>2</sub>, which is comparable to the dihedral angle in Ag<sub>2</sub>O<sub>3</sub> (174.5°). In contrast to the Ag/CeO<sub>2</sub>  
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6 case, where one reduced Ce<sup>3+</sup> ion was observed, here all Ce ions have oxidation state +4 (i.e. all *f*  
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8 states are empty).  
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12 The optimized structures of Cu:H-CeO<sub>2</sub> and Au:H-CeO<sub>2</sub> are shown in Fig. 1b and 1h,  
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14 respectively. Similar to Ag, Au does not introduce major distortions of the surface structure: it  
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16 remains at the same height of the Ce atoms with Au-O bond lengths of 2.12 Å, in the range of the  
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18 observed Au-O bond lengths in Au<sub>2</sub>O<sub>3</sub><sup>33</sup>. In contrast, the Cu atom moves down toward a subsurface  
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20 site with Cu-O bond lengths of 1.89 Å, similar to the observed bond distances in Cu<sub>2</sub>O<sup>34</sup>. The PDOS  
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22 of Cu:H-CeO<sub>2</sub> (Fig. 2b) and Au:H-CeO<sub>2</sub> (Fig. 2h) show metal d states above the Fermi energy, as  
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24 found for Ag (Fig. 2e). The changes of Cu and Au's Bader charges with respect to the adsorbed atoms  
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26 are also similar to those found for Ag. All results are consistent with an oxidation state +3 in the case  
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28 of Au. The picture is less clear for Cu, for which a +2 oxidation state was actually suggested in a  
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30 computational study by Yang et al.<sup>35</sup>  
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### 38 **III.1.3. *M*:CeO<sub>2-x</sub>**

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40 When a surface Ce atom is replaced by Ag, the energy cost to form an oxygen vacancy is  
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42 found to become negative, -0.69 eV. This implies that O-vacancy formation occurs spontaneously in  
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44 the presence of substitutional Ag, as computationally predicted by Shapovalov et al.<sup>36</sup>. For  
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46 comparison, our computed formation energy of a surface oxygen vacancy on metal-free CeO<sub>2</sub>(111)  
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48 surface is 2.05 eV, a value similar to those reported in previous theoretical studies<sup>5 37</sup>. The  
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50 configuration generated by removing one of the surface oxygen atoms bonded to the substitutional  
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52 Ag is shown in Fig. 1f. The excess electrons that result from the creation of the vacancy are acquired  
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54 by the Ag atom. Analysis of the DOS shows indeed that the spin up empty Ag d state of Ag:H-CeO<sub>2</sub>  
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56 (Fig. 2e) becomes occupied in Ag:CeO<sub>2-x</sub> (Fig. 2f), whereas the spin down d state remains empty  
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(dashed lines). This indicates that the Ag oxidation state is +2 in Ag:CeO<sub>2-x</sub>. A similar behavior was reported in computational studies by Szabova et al.<sup>38</sup> and Yang et al.<sup>35</sup> for Cu.

In agreement with Yang et al.<sup>35</sup> and Shapovalov et al.<sup>36</sup>, our calculations predict that the creation of an oxygen vacancy is thermodynamically favored also in the presence of a substitutional Cu, even though the energy gain ( $E_{\text{form}} = -0.12$  eV) is smaller than in the Ag case. In contrast, the creation of an oxygen vacancy is computed to be slightly endothermic,  $E_{\text{form}} = +0.23$  eV, in the presence of substitutional Au. Our result is in agreement with the theoretical value obtained by Kim et al.<sup>39</sup>, whereas Shapovalov et al.<sup>36</sup> predicted a negative formation energy (-0.36 eV) also in the Au case. The PDOS of Cu:CeO<sub>2-x</sub> and Au:CeO<sub>2-x</sub> (Fig. 2c and 2i) are similar to that of Ag:CeO<sub>2-x</sub>, indicating that both excess electrons from the vacancy are acquired by the noble metal. A different result was reported in a theoretical work by Camellone et al.<sup>30</sup>, who found that one of the excess electrons of the oxygen vacancy in Au:CeO<sub>2-x</sub> is acquired by a Ce atom, leading to a reduced Ce<sup>3+</sup>.

### III.2. Pathways of H<sub>2</sub> dissociation

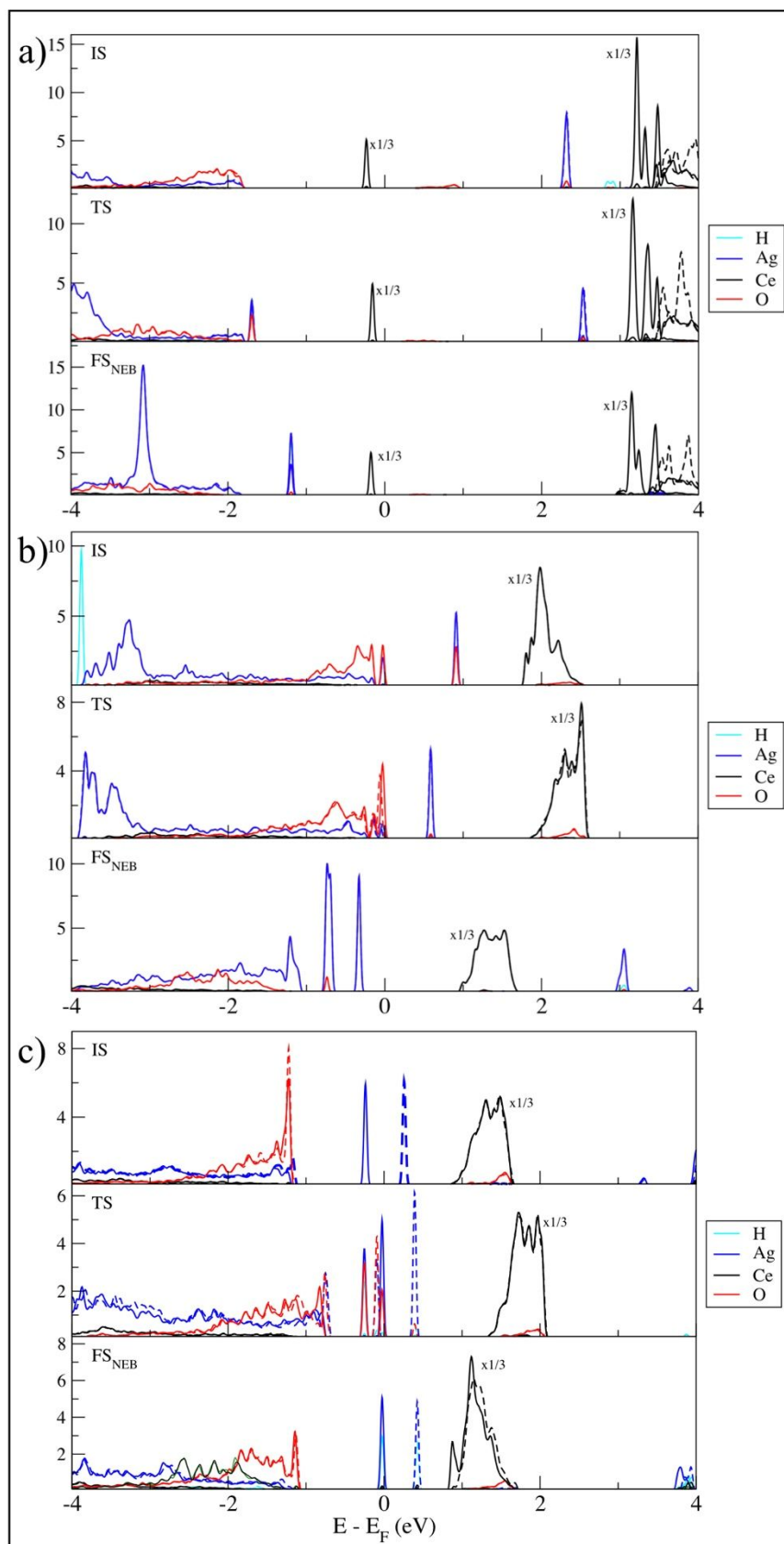
Recent computational studies of H<sub>2</sub> dissociation on metal-free CeO<sub>2</sub> (111) have predicted an heterolytic mechanism leading to the formation of two adjacent OH and the reduction of two surface Ce<sup>4+</sup> cations to Ce<sup>3+</sup><sup>13 40 41</sup>. Following those studies, we calculated the pathways of H<sub>2</sub> dissociation on all the model systems shown in Fig. 1 using the NEB method<sup>25</sup>. In each case, we considered an initial state (IS) where the H<sub>2</sub> molecule is physically adsorbed on the surface and a final state (denoted FS<sub>NEB</sub>) where the molecule is dissociated. For the IS, we always examined several possible configurations of H<sub>2</sub> on the relaxed surface and chose the one energetically most stable. Considering that H<sub>2</sub> dissociation is likely to occur at the metal site, for the FS<sub>NEB</sub> we chose a configuration with one H atom of H<sub>2</sub> bonded to the noble metal atom and the other bonded to a surface oxygen atom, rather than the most stable configuration (denoted FS) with both H atoms adsorbed on surface oxygen atoms. FS<sub>NEB</sub> was found however to evolve spontaneously to FS in the case M:H-CeO<sub>2</sub>.

For better clarity, in the following we will report two reaction energies: the first one,  $\Delta E_{\text{NEB}}$ , is the energy difference between the final and initial states of the NEB calculation; the other,  $\Delta E$ , is the energy of the “true” (most stable) final state with the two H atoms on surface oxygen atoms, referred to the non-interacting gas phase  $\text{H}_2$  molecule and M-CeO<sub>2</sub> system.

### III.2.1. $\text{H}_2$ dissociation on Ag/CeO<sub>2</sub>

The minimum energy pathway (MEP) for  $\text{H}_2$  dissociation on Ag/CeO<sub>2</sub>(111) is shown in Fig. 3. In the IS, the  $\text{H}_2$  molecule is weakly adsorbed on the Ag adatom, with a binding energy of 0.27 eV. The  $\text{H}_2$  bond length is 0.80 Å (vs 0.75 Å in the gas phase) and the Ag-H distance is 1.91 Å. The PDOS shows no charge transfer from  $\text{H}_2$  to the surface (Fig. 4a IS). As final state  $\text{FS}_{\text{NEB}}$ , we consider a configuration with one H bonded to Ag (bond length  $\sim 1.62$  Å), the other bonded to a surface oxygen  $\text{O}_s$  (bond length  $\sim 0.99$  Å), and the Ag adatom bonded to only one surface oxygen (bond length  $\sim 2.15$  Å). The H atom adsorbed onto  $\text{O}_s$  transfers its electron to a surface Ce atom reducing it to  $\text{Ce}^{3+}$ , while the other H reduces the Ag adatom from  $\text{Ag}^+$  to  $\text{Ag}^0$ , as shown by the PDOS in the bottom panel of Fig. 4a, where the Ag s peak is below the Fermi energy. From the MEP in Fig. 3, we can also see that in the initial stages of the reaction the Ag atom rotates towards  $\text{O}_s$ , and breaks one of its O-Ag bonds. At the transition state (TS), the bond length of the  $\text{H}_2$  molecule has increased to 0.90 Å, and one H atom forms a hydrogen bond with  $\text{O}_s$  (dashed line in Fig. 3), with a  $\text{O}_s\text{-H}$  bond length of 1.64 Å, while the other H remains bonded to the Ag atom. The Ag s orbital remains empty and a  $\text{Ce}^{3+}$  peak is still present in the PDOS (Fig. 4a, TS). This TS is only 0.23 eV higher in energy than the IS. We have also determined the barrier for H transfer from Ag ( $\text{FS}_{\text{NEB}}$ ) to a near surface Oxygen (FS), and found it to be 1.0 eV. This barrier is quite high in comparison to that for  $\text{H}_2$  dissociation, suggesting that  $\text{FS}_{\text{NEB}}$  is an observable intermediate of the reaction.





**Fig. 4.** Projected Density of States for the IS, TS, and  $FS_{NEB}$  of  $H_2$  dissociation on: a) Ag/CeO<sub>2</sub>; b) Ag:H-CeO<sub>2</sub>; c) Ag:CeO<sub>2-x</sub>. Full and dashed lines represent the spin up and spin down contributions, respectively.

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3 Similar mechanisms of H<sub>2</sub> dissociation are found on Cu/CeO<sub>2</sub> and Au/CeO<sub>2</sub>, see Figs. S3, S4,  
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5 S5 and S6 of the Supporting Information (SI). In the IS, H<sub>2</sub> is adsorbed on the metal atom with  
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7 adsorption energies of -0.35 and -0.93 eV, for Cu and Au, respectively, and the molecule is more  
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9 stretched than in the Ag case. On Au/CeO<sub>2</sub>, the Au atom shifts from the bridge to the on-top site upon  
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11 H<sub>2</sub> adsorption; a similar change of metal adsorption site was observed also in a theoretical study of  
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13 CO oxidation on Au/CeO<sub>2</sub> by Camellone et al.<sup>30</sup> The existence of a bond between Au and H was  
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15 assumed also by Jaurez et al.<sup>43</sup> to explain their experimental results on H<sub>2</sub> dissociation on Au/CeO<sub>2</sub>.  
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19 The activation energy ( $E_a$ ) and reaction energies ( $\Delta E_{\text{NEB}}$  and  $\Delta E$ ) of H<sub>2</sub> dissociation on the  
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21 three M/CeO<sub>2</sub> surfaces are compared in Table 2, which shows also results for metal-free ceria from  
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23 previous DFT studies. The presence of the noble metal atom on the CeO<sub>2</sub> surface decreases the  
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25 barrier of H<sub>2</sub> dissociation (from 0.99 eV to ~ 0.30 eV) and at the same time makes the reaction energy  
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27 less negative, indicating that a single noble metal atom adsorbed on the CeO<sub>2</sub> (111) is an efficient  
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29 catalyst for H<sub>2</sub> dissociation. The charge transfer from the metal adatom to the CeO<sub>2</sub> surface is crucial  
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31 in favoring the initial binding of the H<sub>2</sub> molecule to the catalyst. This initial binding facilitates the  
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33 dissociation with respect to the case of pure ceria where H<sub>2</sub> does not bind to the surface. The smaller  
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35 absolute value of  $\Delta E$  on M/CeO<sub>2</sub> relative to pure CeO<sub>2</sub> indicates that the H atoms are less strongly  
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37 bound to the surface, which should facilitate the subsequent proton release during PEMFC operation.  
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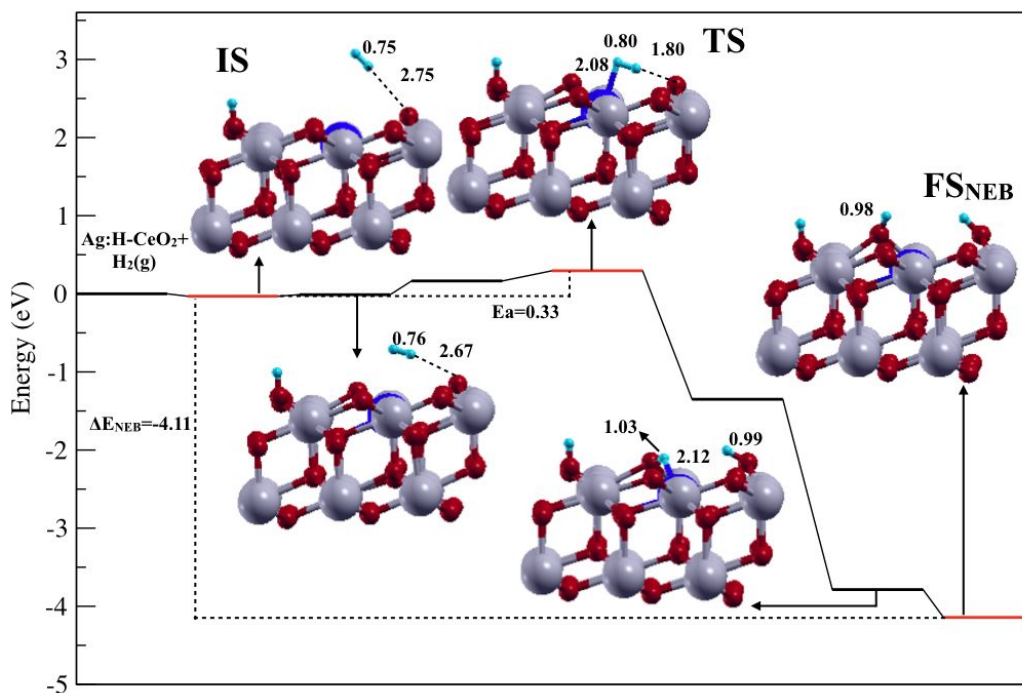
58 **Table 2** Activation energies ( $E_a$ ) and reaction energies ( $\Delta E_{\text{NEB}}$  and  $\Delta E$ ) for H<sub>2</sub> dissociation on pure CeO<sub>2</sub>  
59 (from previous computational studies, as indicated), Cu/CeO<sub>2</sub>, Ag/CeO<sub>2</sub>, and Au/CeO<sub>2</sub> (from this work).  
60  $\Delta E_{\text{NEB}}$  is the energy difference between the final and initial states of the NEB calculation, whereas  $\Delta E$  is the

energy of the H<sub>2</sub> dissociated state with both H atoms on surface oxygens relative to the non-interacting H<sub>2</sub> molecule and M/CeO<sub>2</sub> surface. All energies in eV.

	CeO <sub>2</sub>	Cu	Ag	Au
E <sub>a</sub>	0.99 <sup>13</sup> 0.99 <sup>41</sup> 1.08 <sup>40</sup>	0.29	0.23	0.35
ΔE <sub>NEB</sub>	-2.50 <sup>13</sup> -2.34 <sup>41</sup> -2.50 <sup>40</sup>	-0.55	-0.72	-1.27
ΔE	-2.50 <sup>13</sup> -2.34 <sup>41</sup> -2.50 <sup>40</sup>	-1.99	-1.53	-2.32

### III.2.2. H<sub>2</sub> dissociation on Ag:H-CeO<sub>2</sub>

The MEP for H<sub>2</sub> dissociation on Ag-doped CeO<sub>2</sub> in the presence of an adsorbed H is shown in Fig. 5. In the initial state, H<sub>2</sub> is physically adsorbed (E<sub>ads</sub> = -0.03 eV) with one of its H atoms at distance ~ 2.75 Å from a surface oxygen, and the Ag oxidation state is +3 (Fig. 4b IS). As H<sub>2</sub> moves towards the surface, one of the Ag-O bonds with the surface breaks. At the TS, the H<sub>2</sub> molecule (bond length ~ 0.80 Å) is adsorbed on the Ag atom (Ag-H distance ~ 2.15 Å) while forming also a H-bond with a surface O atom (O<sub>s</sub>) at distance ~ 1.80 Å. From the PDOS it appears that the Ag oxidation state is still +3 at the TS (Fig. 4b TS). After the H-H bond breaks, one hydrogen atom moves directly to O<sub>s</sub> while the other moves spontaneously from Ag to another nearby surface oxygen atom. The activation energy for dissociation, E<sub>a</sub> = 0.33 eV, is significantly lower than in pure CeO<sub>2</sub>. In this system the final state FS<sub>NEB</sub> transforms spontaneously to (i.e. coincides with) FS, H<sub>2</sub> is dissociated with both H atoms adsorbed on surface oxygen atoms and Ag is reduced. As shown by the PDOS in the bottom panel of Fig. 4b, the Ag d states are indeed below E<sub>F</sub>, while the CeO<sub>2</sub> surface remains oxidized, indicating that the Ag oxidation state is +1.



**Fig. 5** Potential energy profile along the MEP of  $\text{H}_2$  dissociation on  $\text{Ag:H-CeO}_2(111)$ . Relevant distances (in  $\text{\AA}$ ) are reported. The color code is the same as in Fig. 3.

The analogous pathways on  $\text{Cu:H-CeO}_2$  and  $\text{Au:H-CeO}_2$  are qualitatively similar to those for  $\text{Ag:H-CeO}_2$  (Figs. S7 and S8), but the charge transfers upon  $\text{H}_2$  dissociation are different. In the Au case, the electrons of both H atoms are transferred not to the noble metal but to the surface, leading to the reduction of two Ce cations from  $\text{Ce}^{4+}$  to  $\text{Ce}^{3+}$ , while Au maintains an oxidation state +3 (Fig. S4b). On  $\text{Cu:H-CeO}_2$ , Cu is reduced to an oxidation state +2 after  $\text{H}_2$  dissociation, as indicated by the PDOS in Fig. S3b.

In Table 3 we compare the activation and reaction energies of  $\text{H}_2$  dissociation on Ag, Au and Cu doped  $\text{CeO}_2$  ( $\text{M:H-CeO}_2$ ). Both the activation and reaction energies are much lower on the metal doped surfaces than on pristine  $\text{CeO}_2(111)$ , and the decrease is more significant for Ag than for Au and Cu.

Charge transfers appear to have a key role in the reaction: Ag facilitates the dissociation of  $\text{H}_2$  by becoming reduced from  $\text{Ag}^{3+}$  to  $\text{Ag}^+$ , its preferred oxidation state, whereas Au is already in its favorite oxidation state (+3) in the IS, so that its effect on  $\text{H}_2$  dissociation is less pronounced. On the other

hand, the strong stability of the final state of H<sub>2</sub> dissociation on M:H-CeO<sub>2</sub>, especially in the Ag case, implies that it is difficult to break the bonds of the adsorbed H atoms with the surface and regenerate the catalyst as required in PEMFCs.

**Table 3** Activation energies ( $E_a$ ) and reaction energies ( $\Delta E_{\text{NEB}}$  and  $\Delta E$ ) of H<sub>2</sub> dissociation on noble metal doped CeO<sub>2</sub> in the presence of adsorbed H (M:H-CeO<sub>2</sub>) or an oxygen vacancy (M:CeO<sub>2-x</sub>).  $\Delta E_{\text{NEB}}$  is the energy difference between the final and initial states of the NEB calculation, while  $\Delta E$  is the energy of the H<sub>2</sub> dissociated state with both H atoms on surface oxygens relative to the non-interacting gas-phase H<sub>2</sub> molecule and doped surface. All energies in eV.

	M:H-CeO <sub>2</sub>			M:CeO <sub>2-x</sub>		
	Cu	Ag	Au	Cu	Ag	Au
$E_a$	0.51	0.33	0.62	0.80	0.34	0.19
$\Delta E_{\text{NEB}}$	-2.36	-4.11	-3.00	-0.55	-0.89	-1.43
$\Delta E$	-2.37	-4.14	-3.03	-2.13	-2.83	-2.89

### III.2.3. H<sub>2</sub> dissociation on Ag:CeO<sub>2-x</sub>

As discussed in Sec. III.1.3, the substitution of a Ce atom by a noble metal favors the formation of oxygen vacancies, so that M:CeO<sub>2-x</sub> is the actual stable configuration of M-doped CeO<sub>2</sub>. The pathway and activation energy for H<sub>2</sub> dissociation on Ag:CeO<sub>2-x</sub> (Fig. 6 and Table 3) are similar to those on Ag:H-CeO<sub>2</sub> (Fig. 5). The initial Ag oxidation state on Ag:CeO<sub>2-x</sub> is +2 (Fig. 4c, IS), and the PDOS changes after H<sub>2</sub> dissociation (Fig. 4c, FS<sub>NEB</sub>) suggest partial transfer of the hydrogen electrons to both Ag and the CeO<sub>2</sub> surface. This is confirmed by the Bader and Löwdin charges, which show an increase of the occupation of the Ag *s* orbital with respect to the IS. However, a clear Ag oxidation state of +1 is observed only after both H atoms are bonded to surface Oxygen atoms. The computed barrier for H diffusion from Ag to a surface Oxygen atom is 0.16 eV, a value significantly lower than the activation energy, 0.34 eV, of H<sub>2</sub> dissociation.

The activation and reaction energies of H<sub>2</sub> dissociation on Cu:CeO<sub>2-x</sub> and Au:CeO<sub>2-x</sub> are summarized in Table 3, while pathways and PDOS are reported in Figs. S9-S10 and S5-S6,

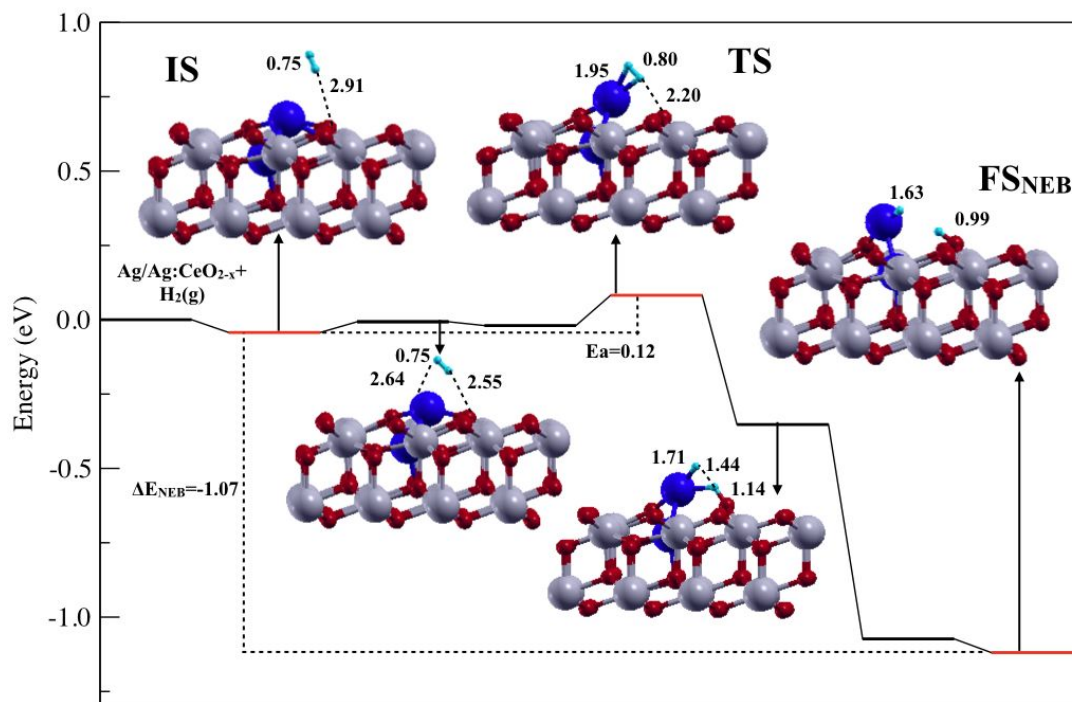
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3 respectively. The barrier is low (0.19 eV) on Au:CeO<sub>2-x</sub>, whereas it is quite high (0.80 eV) on  
4 Cu:CeO<sub>2-x</sub>. A possible reason of this difference is that the initial +2 oxidation state is very stable for  
5 Cu whereas it is not as favorable for Ag and Au, which prefer to take an oxidation state +1, hence  
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7 facilitating the dissociation of the H<sub>2</sub> molecule. On the other hand, the smaller H<sub>2</sub> dissociation barrier  
8 on Au:CeO<sub>2-x</sub> in comparison to Ag:CeO<sub>2-x</sub> may be related to the different TS geometries: on  
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10 Au:CeO<sub>2-x</sub> the H<sub>2</sub> bond at the TS is only slightly longer than at equilibrium (0.78 Å), whereas on  
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12 Ag:CeO<sub>2-x</sub> the H-H bond at the TS is stretched to 0.91 Å, thus requiring more energy.  
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19 The low value of E<sub>a</sub> predicted by our calculations for Au:CeO<sub>2-x</sub> appears to disagree with a  
20 recent study by Qiao et al.<sup>14</sup> showing that atomically dispersed Au on CeO<sub>2</sub> does not activate H<sub>2</sub> at  
21 temperatures below 100 C. On the other hand, our value E<sub>a</sub> = 0.62 eV for H<sub>2</sub> dissociation on H-  
22 Au:CeO<sub>2</sub> is both consistent with the experimental observation by Qiao et al.<sup>14</sup> and also agrees with  
23 their computed activation barrier of 0.69 eV for H<sub>2</sub> dissociation on atomically dispersed Au atoms at  
24 Ce sites (with no oxygen vacancy). Therefore, a possible explanation for the discrepancy of our results  
25 for Au:CeO<sub>2-x</sub> with experiment<sup>14</sup> is the excessive reduction of the Au:CeO<sub>2-x</sub> model, as suggested also  
26 by the positive value of the corresponding O-vacancy formation energy in Sec. III.1.3.  
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37 Table 3 also shows that the reaction energies ΔE are in absolute value smaller on M:CeO<sub>2-x</sub> than  
38 on M:H-CeO<sub>2</sub>, suggesting that the presence of oxygen vacancies improves the activity of noble metal  
39 doped CeO<sub>2</sub> as anode material for PEMFCs.  
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dissociation, one H atom moves onto the surface oxygen to which it was H-bonded while the other stays on the Ag atom. The activation energy of the process, only 0.12 eV, is significantly lower than on pure ceria, as well as on Ag:H-CeO<sub>2</sub> and Ag:CeO<sub>2-x</sub>.



**Fig. 7.** Potential energy profile along the Minimum energy Pathway (MEP) of H<sub>2</sub> dissociation on Ag on Ag:CeO<sub>2-x</sub>(111) as obtained by CI-NEB calculations with seven images. The color code is the same as in Figure 3. Reported distances are in Angstrom.

## CONCLUSIONS

In this work, we have used ab-initio calculations to investigate the mechanism of H<sub>2</sub> dissociation on single Ag, Au and Cu atoms deposited on or doped into CeO<sub>2</sub> (111). An important goal has been to identify an efficient hydrogen oxidation catalyst that might possibly replace Pt as anode material for PEMFCs. We have focused mainly on Ag, whose activity toward H<sub>2</sub> dissociation has been explored for the first time in this work, and considered different surface models, which has allowed us to investigate the effect of different initial oxidation states of the noble metal atom on the activation of H<sub>2</sub>. Our results show that noble metal atoms generally induce a significant decrease of the H<sub>2</sub>

dissociation barrier with respect to the pure ceria surface, and the charge transfers involved in the reaction play a crucial role in this decrease. For noble metal adatoms on pristine ceria ( $M/CeO_2$ ), the initial metal oxidation state is +1 and the charge transfers involved in  $H_2$  dissociation are the same for all metals.  $Ag/CeO_2$  shows the lowest activation energy as well as the less negative reaction energy among the different supported metals (see Table 2), which makes it a potentially interesting material for PEMFC anodes. In the case where the noble metal is substitutional to a surface Ce atom, the formation of oxygen vacancies is highly favored ( $M:CeO_{2-x}$ ). The initial metal oxidation state is +2, and among the different metals Au has the lowest activation energy, closely followed by Ag, which also shows a slightly less negative reaction energy (Table 3). Finally, as an alternative to the formation of an oxygen vacancy, we also considered the case where the missing charge of the substitutional noble metal atom is partially compensated by an adsorbed H ( $M:H-CeO_2$ ). In this case the initial metal oxidation state is +3 and Ag induces the largest decrease of the activation energy with respect to pure ceria (Table 3). Ag plays an active role by acquiring the electrons of the dissociated  $H_2$  molecule, so that its oxidation state changes from +3 to +1, which is the optimal one for Ag.

In general, we found that when the noble metal atom is in its favored oxidation state in the initial state of the reaction, the energy barrier lowering with respect to pure ceria is small: this is the case of  $H_2$  dissociation on  $Au:H-CeO_2$  or  $Cu:CeO_{2-x}$ . Except for these cases, the noble metal participates actively in the reaction and contributes to the barrier lowering by accepting electrons from  $H_2$  so as to reach its favorite oxidation state. Finally, our results indicate that  $CeO_2$  modified with single Ag atoms is a promising candidate material for replacing Platinum in the anode of PEMFCs.

## ASSOCIATED CONTENT

### Supporting Information

Top view of the different configurations (Fig. S1), potential energy profile of  $H_2$  dissociation for  $Ag/CeO_2$  using  $U=4.5$  eV (Fig. S2), potential energy profile of  $H_2$  dissociation for  $Cu/CeO_2$  and  $Au/CeO_2$  (Fig. S3 and S4), projected density of states (Fig. S5 and S6), potential energy profile of

H<sub>2</sub> dissociation for Cu:H-CeO<sub>2</sub> and Au:H-CeO<sub>2</sub> (Fig. S7 and S8), potential energy profile of H<sub>2</sub> dissociation for Cu:CeO<sub>2-x</sub> and Au:CeO<sub>2-x</sub> (Fig. S9 and S10).

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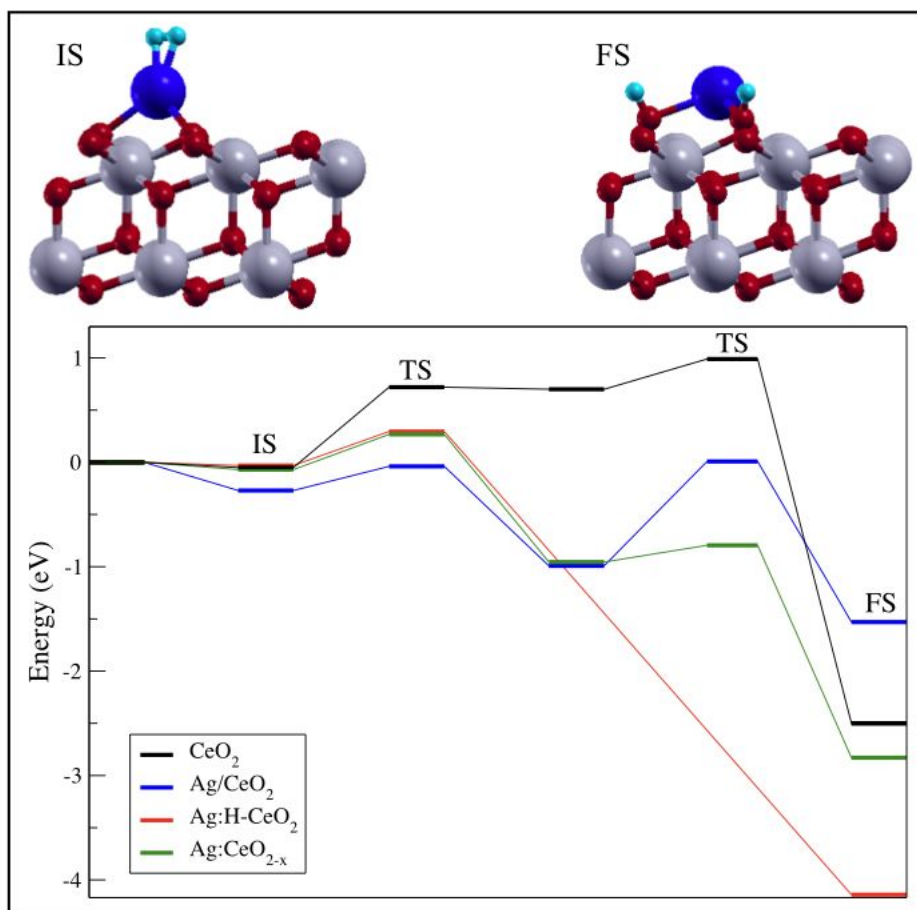


Figure Toc graphic